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us take $z = 4cy^2$ We can write the

Lagrangian Equations for this motion 1

$T = m(\dot{r}^2 + r^2 \dot{\theta}^2 + \dot{z}^2)$ 2 $U = mgz$

In our case $r = y$ and $z = cy^2$ so we

can say that $\dot{z} = 2yc\dot{y}$ and we know

that $\dot{y} = \dot{y}t$ and $\dot{\theta} = \dot{\theta}$ Now we can

write the Lagrangian as $L = T - U$ 1 L

$= m(\dot{y}^2 + \dots$

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4.1 Part (a) In the given con guration,

both springs elongate or compress by

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~~Solutions~~ the same magnitude. Suppose q denotes the position of the mass m from the left end. At $t=0$, $q(0) = a=2$, but the unstretched lengths of both springs are given to ...

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$1 + k \cdot 2)(Q + b \sin(\omega t))^2$ (54) The
Hamiltonian is now explicitly
dependent on time, and hence is not
conserved, as is confirmed by the fact
that $dH/dt \neq 0$. The energy is given by
 $E = T + V = \frac{1}{2} (\dot{Q} + b \omega \cos(\omega t))^2 + \frac{1}{2}$
 $(k \cdot 1 + k \cdot 2)(Q + b \sin(\omega t))^2$ (55) So, dE/dt

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 $\cong m(Q_1 + b \cos(\theta))(Q_1^2 \sin(\theta)) + (k_1 + k_2) Q_1^2$

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Chapter 7 now presents special relativity using the standard real metric (in lieu of the complex Minkowski space) and

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CHAPTER 9 – CANONICAL TRANSFORMATIONS

DERIVATIONS: 9.4. Show directly that the transformation is canonical. 9.4.

Sol. We are given a transformation as follows, We know that the fundamental Poisson Brackets of the transformed variables have the same value when evaluated with respect to any canonical coordinate set.

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We use the first constraint to solve for the coordinate r : $r = R + a$, $r' = r'' = 0$. We use this solution in Lagrange's equations for r , θ : $m(R + a)\ddot{\theta} + mg \sin \theta = 0$ (6) $m(R + a)\ddot{r} + mg(R + a) \cos \theta = \mu(R + a)$ (7) We use the rolling constraint to find an expression for $\dot{\theta}$ as a function of \dot{r} : $\dot{\theta} = \frac{a + R}{a} \dot{r} + \dot{\theta}_0$ (8)

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